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# THE LANCZOS ALGOR ITHM FOR THE SYMMETRIC $A x=\lambda B x$ PROBLEM 

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## ABSTRACT

The problem of computing the eigensystem of $A x=\lambda B x$ when $A$ and B are symmetric and B is positive definite is considered. A generalization of the Lanczos algorithm for reducing the problem to a symmetric tridiagonal eigenproblem is given. A numerically stable variant of the algorithm is described. The new algorithm depends heavily upon the computation of elementary Hermitian matrices. An AIGOL $W$ procedure and a numerical example are also given.

The Lanczos Algorithm for the Symmetric $A x=\lambda B x$ Problem

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## 1. Theoretical Background

In many fields of work the solution of the eigenproblem

$$
\begin{equation*}
A x=\lambda B x \tag{1}
\end{equation*}
$$

is required where $A$ and $B$ are symmetric and $B$ is positive definite. This problem can be reduced to the standard symmetric eigenproblem by making use of the Cholesky factorization of $B$ defined by

$$
\begin{equation*}
B=L L^{T} \tag{2}
\end{equation*}
$$

Equation (1) is then equivalent to

$$
\begin{equation*}
L^{-1} A L^{-T}\left(I^{T} x\right)=\lambda\left(I^{T} x\right) \tag{3}
\end{equation*}
$$

and $L^{-1} A L^{-T}$ is a real symmetric matrix.
When A and B are narrow symmetric band matrices of high order this reduction has the disadvantage that $L^{-1} A L^{-T}$ is, in general, a full matrix. However, $I$ itself is of band form and hence we can certainly multiply an arbitrary vector by $L^{-1} A L^{-T}$ in an economical manner. In fact, if we write

$$
\begin{equation*}
\mathrm{z}=\mathrm{L}^{-1} A L^{-T} \mathrm{~T} \tag{4}
\end{equation*}
$$

then $z$ can be determined in the steps

$$
\begin{equation*}
L^{T} y=x, \quad A y=w, \quad L z=w \tag{5}
\end{equation*}
$$

and in this way we can take full advantage of the band forms of L and A. The total number of multiplications in the determination of $z$ is only marginally greater than in the determination of both $A x$ and Bx taking advantage of the band forms of $A$ and $B$.

Now in the Lanczos algorithm for a symmetric matrix C , the only way in which $C$ is used is in the pre-multiplication of vectors. The algorithm may be described as follows. Let $x_{I}$ be an arbitrary unit vector $\left(\left\|x_{1}\right\|_{2}=1\right)$; then determine sequences of vectors $y_{r}$ and $x_{r}$ defined by

$$
\begin{align*}
& \mathrm{y}_{2}=C \mathrm{x}_{1}-\alpha_{1} \mathrm{x}_{1}, \gamma_{2} \mathrm{x}_{2}=\mathrm{y}_{2}, \mathrm{I}_{\mathrm{I}} \|_{2}=1, \mathrm{Y} 2>0  \tag{6}\\
& \mathrm{y}_{r+1}=C x_{r}-\alpha_{r} x_{r}-\beta_{r} x_{r-1},{ }_{y_{r+1}} x_{r+1}=y_{r+1}, \\
& \left\|x_{r+1}\right\|_{2}=1, \gamma_{r+1}>0 \tag{7}
\end{align*}
$$

where the sequence is continued until $y_{r+1}=\theta$. The $\alpha_{r}$ and $\beta_{r}$ are determined so that $\mathrm{y}_{\mathrm{r}+1}$ is orthogonal to $\mathrm{x}_{\mathrm{r}}$ and $\mathrm{x}_{\mathrm{r}-1}$ and $\alpha_{1}$ so that $y_{2}$ is orthogonal to $\mathrm{x}_{1}$, These relations ensure that $\gamma_{r}=\beta_{r}$ and $y_{r+1}$ is automatically orthogonal to $x_{1}, x_{2}, \ldots, x_{r}$. Notice that when $y_{r+1}$ has been determined, Equations (6) and (7) imply that

$$
\begin{equation*}
\mathrm{C}\left[\mathrm{x}_{1}, \mathrm{x}_{2}, \ldots, \mathrm{x}_{r}\right]=\left[\mathrm{x}_{1}, \mathrm{x}_{2}, \ldots, \mathrm{x}_{r}\right] \mathrm{T}_{r}+\left[0,0, \ldots, 0, \mathrm{y}_{r+1}\right] \tag{8}
\end{equation*}
$$

where $\mathrm{T}_{\mathrm{r}}$ is the tridiagonal matrix with diagonal elements equal to the $a_{1}$, superdiagonal elements equal to the $\beta_{i}$ and subdiagonal elements equal to the $\gamma_{i}$. This is true even if $\alpha_{i}$ and $\beta_{i}$ are chosen arbitrarily! If $y_{r+1}=\theta$, then (8) gives

$$
\begin{equation*}
C\left[x_{1}, x_{2}, \ldots, x_{r}\right]=\left[x_{1}, x_{2}, \quad, ., x_{r}\right] \mathbb{I}_{r} \tag{9}
\end{equation*}
$$

and provided only the $x_{i}$ are independent, $T_{r}$ gives $r$ of the eigenvalues of $C$ and enables us to compute the corresponding eigenvectors. In the Lanczos algorithm, the orthonormality of the $x_{i}$ ensures their independence and also the symmetry of $T_{r}$.

Since $\bar{y}_{r+l}$ is orthogonal to the $r$ orthogonal vectors $x_{1}, x_{2}, \ldots, x_{r}$, the process must terminate with $\mathrm{y}_{\mathrm{n}+1}$ if it has not done so before. In fact, we may regard the Lanczos algorithm as a method of ensuring that the $x_{i}$ are independent and that the process does terminate. When the Lanczos algorithm terminates before $y_{n+1}$, then we have

$$
\begin{equation*}
0=y_{r+1}=C x_{r}-a_{r}{\underset{r}{x}}^{x_{r}}-\beta_{r} x_{r-1} \tag{10}
\end{equation*}
$$

If we choose $x_{r+1}$ to be any unit vector orthogonal to $x_{1}, \ldots, x_{r}$, then Equation (10) gives

$$
\begin{equation*}
C x_{r}-\alpha_{r} x_{r}-\beta_{r} x_{r-1}=\theta=0 \cdot x_{r+1} \tag{11}
\end{equation*}
$$

Hence we may take $\gamma_{r+1}=0$ and continue the algorithm with $\mathrm{x}_{\mathrm{r}+1}$, This again ensures that subsequent $X_{i}$ are orthogonal to all earlier $x_{j}$. We may restart as often as necessary until we finally reach the null $y_{n+1}$. If the termination takes place after $r_{1}, r_{1}+r_{2}, r_{1}+r_{2}+r_{3}, \ldots, r_{1}+r_{2}+\ldots+r_{k}$ steps then we have

$$
\begin{equation*}
C\left[x_{1}, x_{2}, \ldots, x_{n}\right]=\left[x_{1}, x_{2}, 1\right. \text { 㧙明䡍 } \tag{12}
\end{equation*}
$$

where T is the direct sum of k symmetric tridiagonalmatrices of orders $r_{1}, r_{2}, \boxtimes \Omega \&$ ．The eigensystem of $T$ gives the eigensystem of $c$ ．Premature termination of the sequence should not be regarded as a breakdown；in fact it leads to a simplification since it is easier to find the eigenvalues and eigenvectors of several smaller tridiagonal matrices than of one large one．The only disadvantage is the necessity for determining the restarting vectors．

At first sight the Lanczos algorithm is very attractive whenever $C$ is sparse（including，in particular，the case when $C=L^{-1} A L^{-T}$ ）． Except when restarting，we need only the two vectors $X_{r}$ and $x_{r-1}$ to determine $y_{r+1}$ and $x_{r+1}$ and hence requirements on the high－speed store appear to be very modest．Unfortunately，if the algorithm is carried out as described，the later $\mathrm{x}_{\mathrm{r}}$ may be very far from orthogonal to the earlier ones．When this is true，we have no guarantee that $\mathrm{y}_{\mathrm{n}+1}$ will be null to working accuracy．Moreover，we may get near linear dependency of $x_{1}, x_{2}, \ldots, x_{r}$ for some $r<n+1$ ．

This departure from orthogonality is sometimes said to be the result of accumulation of rounding errors，but this is very misleading． It occurs when there is a good deal of cancellation when computing $y_{r+1}$ from $C x_{r}-\alpha_{r} x_{r}-\beta_{r} x_{r}$ ．This cancellation can occur even in the determination of $y_{2}$ when $C$ is of order 2 （say），and in this case one could scarcely speak of accumulation of rounding errors．Before discussing how to obtain a complete set of orthonormal $x_{i}$ we consider
the implication of Equation (8) when rounding errors are taken into account. We now have

$$
\begin{equation*}
C\left[x_{1}, x_{2}, \ldots, x_{r}\right]=\left[x_{1}, x_{2}, \ldots, x_{r}\right] T_{r}+\left[0,0, \ldots, y_{r+1}\right]-\left[\varepsilon_{1}, \varepsilon_{2}, \ldots, \varepsilon_{r}\right] \tag{13}
\end{equation*}
$$

where $\varepsilon_{i}$ is defined by the relation

$$
\begin{equation*}
\gamma_{i+1} x_{i+1}=C x_{i}-\alpha_{i} x_{j}-\beta_{i-1}+\varepsilon_{i}, \tag{14}
\end{equation*}
$$

the $x_{i+1}$ denoting the computed $x_{i+1}$. If $\lambda$ and $z$ are an eigenvalue and eigenvector of $T_{r}$ we have

$$
\begin{equation*}
T_{r} Z=\lambda z \tag{15}
\end{equation*}
$$

and hence

$$
\begin{equation*}
C\left[x_{1}, x_{2}, \ldots, x_{r}\right] z=\lambda\left[x_{1}, x_{2}, \ldots, x_{r}\right] z+z_{r}{\underset{r}{r}+1}_{-}\left[\varepsilon_{1}, \varepsilon_{2}, \ldots, \varepsilon_{r}\right] z \tag{16}
\end{equation*}
$$

where the $z_{i}$ are the components of $z$. Now $z$ may be taken to be a unit vector and $\left[\varepsilon_{1}, \varepsilon_{2}, \ldots, \varepsilon_{r}\right] z$ is therefore of the order of machepsx\|c\| where macheps is the machine precision. Hence if $Z_{r}{ }^{Y_{r+1}}$ is also of the order of machepsx\|c\| we have

$$
\begin{equation*}
\mathrm{CW}=\lambda_{\mathrm{w}}+\mathrm{e} \tag{17}
\end{equation*}
$$

where $w=\left[x_{1}, x_{2}, \ldots, x_{r}\right] z$ and $\|e\|$ is of order machepsx|C\|. This shows that if we reach a $y_{r+1}$ which is negligible to working accuracy, eigenvalues and eigenvectors of $T_{r}$ give good approximations to eigenvalues and eigenvectors of $C$. (The latter must be interpreted in terms of exact eigenvalues of some $C+E$ where $E$ is small.) Also,
provided the $x_{1}, \ldots, x_{r}$ have not yet departed too far from orthogonality, we shall indeed get good approximations to $r$ eigenvalues and $r$ eigenvectors of $C$.

However, even if $Y_{r+1}$ is not itself negligible it may happen that for some eigenvalues $\lambda$ of $T_{r}$, the corresponding $z$ may have a moderately small $z_{r}$. Then $z_{r} y_{r+1}$ may be negligible even if $y_{r+1}$ is not sufficiently small in itself. This situation may be induced to some extent by starting with an initial vector $\mathrm{x}_{1}$ which consists mainly of a linear combination of a few dominant eigenvectors; such an initial vector is obtained if an arbitrary vector $x_{0}$ is premultiplied several times with $C$. It cannot be too strongly emphasized that the size of the vectors $\boldsymbol{\varepsilon}_{i}$ is not in any way affected by cancellation or by the normalization of $y_{i+1}$ to give $x_{i+1}$, the vector $\varepsilon_{i}$ consists entirely of the rounding errors made in actually multiplying the computed $x_{i}$ by $C$ and subtracting multiples of $x_{i}$ and $x_{i-1}$. Even the accuracy of the $\alpha_{i}$ and $\beta_{1}$ is quite irrelevant in so far as it effects the size of the $\varepsilon_{i}$ though if they were chosen at random it is unlikely that a small $y_{r+1}$ would emerge. These considerations show - why the Lanczos algorithm often gives remarkably accurate approximations to dominant eigenvalues and eigenvectors after quite a few steps.

In order to be certain of obtaining the full set of eigenvalues and eigenvectors it is necessary to ensure that the computed $\mathrm{x}_{\mathrm{i}}$ are orthogonal to working accuracy. The conventional way of doing this is as follows. After computing $\mathrm{y}_{\mathrm{r}+1}$ via Equations (7), it is reorthogonalized with respect to $x_{r}, x_{r-1}, x_{1}$ (N.B. Since the lack of orthogonality
is caused by cancellation and not by the accumulation of rounding errors, it is just as necessary to reorthogonalize with respect to $\mathrm{X}_{\mathrm{r}}$ and $\mathbf{x}_{r-1}$ (although $y_{r+1}$ has just be orthogonalized with respect to these vectors) as to the earlier vectors.) We may write

$$
\begin{equation*}
y_{r+1}=y_{r+1}-\eta_{1} x_{1}-\eta_{2} x_{2}-\cdots-\eta_{r} x_{r} \tag{18}
\end{equation*}
$$

where the $\eta_{i}$ are chosen so as to give orthogonality. If cancellation takes place when deriving the $\overline{\mathrm{y}}_{\mathrm{r}+1}$, then $\overline{\mathrm{y}}_{\mathrm{r}+1}$ must be reorthogonalized yet again with respect to all earlier vectors. Moreover, if $y_{r+1}$ or $\mathrm{Y}_{\mathrm{r}+1}$ vanishes, then a technique is needed for restarting. Although a perfectly satisfactory procedure may be constructed on these lines, it is not aesthetically pleasing. We now describe an alternative procedure for ensuring the orthogonality of the computed $x_{i}$ to working accuracy.

Suppose $x_{1}, x_{2}, \ldots, x_{r}$ have already been determined and are orthogonal to working accuracy, and that the matrix $X_{r}=\left[X_{1}, X_{2}, \ldots \phi X_{r}\right]$ has been reduced to upper-triangular form by premultiplication with $r$ elementary Hermitian matrices $P_{1}, P_{2}, \ldots, P_{r}$. Here

$$
\begin{equation*}
P_{i}=I-2 w_{i} w_{i}^{T},\left\|w_{i}\right\|_{2}=1 \tag{19}
\end{equation*}
$$

and the first i-l components of $w_{i}$ are zero. From the orthonormality of the $x_{i}$, this reduced form must consist of the first $r$ columns of $I$. The vector $\mathrm{y}_{\mathrm{r}+1}$ is then determined from Equation (7) and the vector $\mathrm{z}_{\mathrm{r}+1}$ is determined from the relation

$$
\begin{equation*}
z_{r+1}=P_{r} \cdot \cdot P_{2} P_{1} y_{r+1} . \tag{20}
\end{equation*}
$$

Now with exact computation $\mathrm{y}_{\mathrm{r}+1}$ would be orthogonal to $\mathrm{x}_{1}, \ldots, \mathrm{x}_{\mathrm{r}}$ and hence ${ }^{Z}{ }_{r+1}$ would be orthogonal to the $r$ vectors $P_{r} . . P_{I} x_{i}$, $(i=I, \ldots, r)$. But these vectors are $e_{1}, e_{2}, \ldots, e_{r}$, and hence with exact computations $z_{r+1}$ would have zero components in elements 1 to $r$. We now determine $P_{r+1}$ so that $P_{r+1}^{Z_{r+1}}$ has zero components in elements $r+2, \ldots, n$. With exact computation $P_{r+1} Z_{r+1}$ would be a multiple of $e_{r+1}$. Now define $x_{r+1}$ by the relation

$$
\begin{equation*}
x_{r+1}=P_{1} P_{2} \ldots P_{r+1} e_{r+1} \tag{21}
\end{equation*}
$$

so that $X_{r+1}$ is automatically a unit vector. If all the computation had been exact $x_{r+1}$ would merely be a multiple of $y_{r+1}$, the multiple being chosen so as to make $x_{r+1}$ a unit vector. Notice that this technique gives us a method of continuing when $y_{r+1}=\theta$. We merely define $x_{r+1}$ by the relation

$$
\begin{equation*}
x_{r+1}=P_{1} P_{2} . . \quad P_{r} e_{r+1} \tag{22}
\end{equation*}
$$

which corresponds to having taken $I$ in place of $P_{r+1}$ in Equation (21). From the derivation of the $P_{i}$, the vector $X_{r+1}$ is automatically orthogonal to $x_{1}, x_{2}, \ldots, x_{r}$ since

$$
\begin{equation*}
\mathbf{x}_{i}=P_{1} P_{2} \cdots P_{i} e_{i}=P_{1} P_{2} \cdots P_{i} P_{i+1} \cdots P_{r} e_{i} \quad(i \leq r) \tag{23}
\end{equation*}
$$

giving

$$
\begin{equation*}
x_{r+1}^{T} X_{i}=e_{r+1}^{T} P_{r}^{-T} \cdots P_{2}^{T} P_{1}^{T} P_{1} P_{2} \ldots P_{r} e_{i}=e_{r+1} e_{i}=0 \tag{24}
\end{equation*}
$$

Notice that with this technique, when determining $y_{r+1}$ we need only $\mathbf{x}_{r}$ and $x_{r-1}$. To determine an $X_{r+1}$ accurately orthogonal to $x_{1}, \ldots, x_{r}$ we need only $P_{1}, \ldots, P_{r}$ and these may be stored via the corresponding $w_{i}$. Since the first $i-l$ components of $w_{i}$ are zero, approximately $\frac{1}{2} n^{2}$ registers are needed to store full information on the $P_{i}$ as against $n^{2}$ if we store the $x_{i}$. The full set of Equations (21) shows that

$$
\begin{equation*}
\left[x_{1}, x_{2}, \ldots, x_{n}\right]=P_{1} P_{2} \ldots P_{n} \tag{25}
\end{equation*}
$$

so that in retaining information on the $P_{i}$ via the $w_{i}$ we effectively have full information on the $\mathrm{x}_{\mathrm{i}}$.

The quantity $\gamma_{i}$ may be derived via the relation

$$
\begin{equation*}
\gamma_{i}=\left\|y_{i-1}\right\|_{2} \tag{26}
\end{equation*}
$$

We can take $\beta_{i}$ to be equal to the $\gamma_{i}$ derived in this way or we can determine it by making $y_{i+1}$ orthogonal to $x_{i-1}$, that is, via the relation

$$
\begin{equation*}
\beta_{i}=x_{i-1}^{T} x_{i} \tag{27}
\end{equation*}
$$

Even after the reorthogonalization $\gamma_{i}$ and $\beta_{i}$ determined in this way will agree to within a small multiple of machepsx $\|C\|_{2}$. In practice, it is instructive to compare the $\beta_{i}$ and $\gamma_{i}$ obtained from Equations (26) and (27). We have finally

$$
\begin{equation*}
C\left[x_{1}, x_{2}, \ldots, x_{n}\right]=\left[x_{1}, x_{2}, \quad 0\right. \text { X } \tag{28}
\end{equation*}
$$

to working accuracy, where the $x_{i}$ are orthogonal to working accuracy and T is a symmetric tridiagonalmatrix. If

$$
\begin{equation*}
T z_{i}=\lambda_{i} z_{i} \tag{29}
\end{equation*}
$$

then
giving

$$
\begin{equation*}
C\left(P_{1} P_{2} \cdots P_{n}\right) z_{i}=\lambda_{i}\left(P_{-i} \cdots P\right) z_{n i}, \quad \text { or } \quad C Q z_{i}=\lambda_{i}^{Q z_{i}} \tag{31}
\end{equation*}
$$

so that the $\widetilde{\mathrm{P}}_{\mathrm{j}}$ give sufficient information to enable us to determine eigenvectors of $C$. When $C=L^{-1} A L^{-T}$, we have

$$
\begin{equation*}
L^{-1} A L^{-T} Q_{i}=\lambda_{i} Q z_{i} \quad \text { or } A\left(L^{-T_{Q z_{i}}}\right)=\lambda_{i} B\left(L^{-T_{Q Z}}\right) \tag{32}
\end{equation*}
$$

and hence eigenvectors of $A-\lambda B$ may be determined using the matrix $L$. If the $z_{i}$ are a set of orthonormal eigenvectors of $T$, then $p_{i}=I^{-T} Q z_{i}$ gives a set of eigenvectors for the problem $A p=\lambda B p$ such that

$$
\begin{align*}
& p_{i}^{T} B p_{i}=z_{i}^{T} Q^{T} L^{-} I_{L L^{\Gamma}} \bar{L}^{T} Q z_{i}=z_{i}^{T} z_{i}=1  \tag{33}\\
& p_{i}^{T} B p_{j}=z_{i}^{T} Q^{T} L^{-1} L L^{T} L^{-T} Q z_{j}=z_{i}^{T} z_{j}=0 \quad(i \neq j) \tag{34}
\end{align*}
$$

## 2. Applicability

reducb may be used to reduce the eigenproblem $A x=\lambda B x$ to the standard symmetric eigenproblem $T y=\lambda y$ where $T$ is tridiagonal.

While reducb may be used whenever $A$ and $B$ are symmetric and positive definite, it is best used in problems in which the band width of $A$ and $B$ are small in comparison to their order.

The derived tridiagonal system may be solved by a variety of methods [4]. The eigenvalues of the derived standard problem are those of the original problem, but the vectors are related as indicated by Equations (31) and (32).

## 3. Formal Parameter List

Input to procedure reducb.
n order of matrices $A$ and $B$.
ma number of lower diagonals of A.
mb number of lower diagonals of $B$.
a elements of the lower triangle of the symmetric matrix $A$ stored as an $n x(m a+1)$ array.
b elements of the symmetric matrix $B$ stored as an $\mathrm{n} \times(\mathrm{mb}+\mathrm{l})$ array.

Output of procedure reducb.
alpha diagonal elements of the symmetric tridiagonalmatrix $T$ similar to $\bar{L}^{1} A \bar{L}^{T}$.
beta codiagonal elements of $T$.
b the lower triangle of $L$ such that $L L^{T}=B$, stored as an $n \times(m b+1)$ array (overwriting the original in $b$ ).

U information on the matrices $P_{i}=I-2 w_{i} W_{i}^{T} \quad$ (This may be stored as an $n \times n$ array, but for economy it can be stored as a linear array of order $\left.\frac{1}{2} n(n+1).\right)$
fail exit used if B , possibly as the result of rounding errors, is not positive definite.

```
procedure reducb (integer value n, ma, mb:
    long real array a, b(*,*);
    long real array alpha, beta(*);
    long real array u(*,*);
    procedure fall):
```

comment Reduction of the symmetric eigenvalue problem
$A x=1 a m b d a * B x$
with symmetric band matrix $A$ and symmetric positive
long real yo.y1,z;
long real array $v, \times 0, \times 1, y(1:: n)$;
comment Compute the Cholesky factor of $B$;
for $:=1$ step 1 until $n$ do
begin $p:=$ (if $1>m b$ then 0 else $m b-i+1)$;
$r:=i-m b+p$;
for $1:=p$ step 1 until mb do
begin $s:=j-1$;
$\mathrm{q}:=\mathrm{mb}-\mathrm{j}+\mathrm{p}$;
z:=b(i,j);
for $k:=p$ step 1 unt il s do
begin $z:=z-b(1, k) * b(r, q)$;
$\mathrm{q}:=\mathrm{q}+1$;
end:
if $j=m b$ then
begin if $z<0$ then fall;
$b(i, j):=$ longsqrt(z);
end
else $b(i, j):=z / b(r, m b)$;
$r:=r+1$;
end $j$;
end formi:

```
comment Compute tridiagonal form;
beta(1):=01; y(1):=11;
for i:=2 step 1 until n do y(i):=01;
for k:=1 step 1 unt1] n do
begin
    z:=01;
    for j:=1 step 1 until n do
    begin x0(j):=x1(j);
        xl(j):=y(j);
        z:=z+y(j)*y(j);
    end;
    y0:=y1; y1:=z;
    comment Multiply x by inv(L)*A*inv(L');
    s:=mb-1:
    for i:=n step - 1 until 1 do
    begin p:=(if i< =n-mb then 0 else mb+i-n);
        q:=i;
        z:=xl(1);
        forj:=s step -1 until p do
        begin a:=q+1;
        z:=z-b(q,j)*v(q);
        end:
        v(i):=z/b(i,mb);
    end lve;
    for i:=1 step1until n do
    begin p:=(lf i>ma then 0 else ma-i+1);
        q:=i-ma+p;
        z:=01;
        for i:=p stepluntil ma do
        begin z:=z+a(i,j)*v(q):
        q: =q+1;
        end;
        p:=(if i<=n-ma then 0 else ma+i-n);
        f o rj :=ma-1 sted -1 untilp do
        begin z:=z+a(q,j)*v(q);
        q:=q+1;
        end;
        y(i):=z;
    end av:
    for i:=1 step 1 until n do
    begin p:=(lf i>mb then 0 else mb-i+1);
    q:=1;
    z:=y(i);
    for j:=s step -1 until p do
        begin q:=q-1;
        z:=z-b(i,j)*v(a):
        end;
        v(i):=z/b(i,mb);
    end solve;
```

```
    comment Compute alpha(k) andbeta(k);
z:=01;
for j:=1 stepluntil n do
z:=z+xl(j)*v(j);
alpha(k):=z:=z/y1;
for j:=1 step1 unt il n do
y(j):=v(j)-z*xl(j);
if }\mp@subsup{k}{}{-}=1\mathrm{ then
begin z:=01;
    for j:=1 stepluntil n do
    z:=z+x0(j)*y(j);
    beta(k):=z:=z/y0;
    if k=n then go to 11;
    for j:=1 stepl until n do
    y(j):=y(j)-z*x0(j);
end:
```

comment Normal ize and reorthogonal ize $y$ with respect to previous col umns of $X$;
for i:=2 step 1 until $k$ do
begon:
$\frac{\text { for }}{z:=z / u(i)} \frac{\text { step } 1}{1)}$ until $n$ do $z:=z+u(j, i) * y(j)$;
for $j:=i$ step 1 until $n$ do $y(j):=y(j)-z * u(j, i)$;
end:
z:=01;
for $1:=k+1$ stepluntil $n$ do $z:=z+y(i) * y(i)$;
if $z=01 \mathrm{then}$
begin, 1): $=11$;
for $i:=k+1$ step 1 until $n$ do $u(i, k+1):=01$; end else
begin
$z:=$ if $y(k+1)\rangle=01$ then. longsqrt $(z)$ else -longsqrt $(z)$;
$u(k+1, k+1):=y(k+1)+z ; u(k+1,1):=u(k+1, k+1) * z ;$
for $j:=k+2$ step 1 until $n$ do $u(j, k+1):=y(j)$;
end:
for $j:=1$ step 1 until $n$ do $y(j):=$ if $j=k+1$ then 11 else 01 ;
for i : $=k+1$ step -1 until 2 do
begin $z:=01$
for $j:=i$ step 1 until $n$ do $z:=z+u(j, i) * y(j) ;$ z:=z/u(i,1);
for $j:=i$ step 1 until $n$ do $y(j):=y(j)-z * u(j, i)$; end:

11:
end $k$;
end reducb:
5. Organizational and Notational Details

The lower triangle of $A$ is stored in such a manner that array element $a(i, m a-i+j)$ contains the value of matrix element $A(i, j)$, $i=1, . \operatorname{n}$, and $j=\max (i-m a, 1), \ldots i$. Thus, columns of a correspond to diagonals of A . B is stored similarly. L , the Cholesky factor of $B$, is lower triangular with the same number of diagonals as the lower triangle of $B$. L is stored as. B is, overwriting B in b .

The initial vector $x_{1}$ is chosen to be $e_{1}$. The details of the elementary Hermitian matrices $P_{1}$ are contained in the vectors $U i$ and the scalars $K_{i}$, where

$$
P_{i}=I-\frac{U_{i}^{U}{ }_{i}^{t}}{K_{i}}
$$

If it should happen that $z_{i}=\theta$, cf. (20), then $U_{i}=\theta$, and $K_{i}=\mathbf{1}$, so that the corresponding $P_{i}$ is the identity matrix. Also, since $x_{1}$ is $e_{1}$, $P_{1}$ is chosen to be the identity matrix, and information on $P_{1}$ is not stored. Otherwise, Ui is stored in the ith column of $u$, and $K_{i}$ is stored in $u(i, 1), i=2, \ldots, n$.

The diagonal of the reduced symmetric tridiagonal matrix is stored in the array alpha, and the off-diagonal in the last $n-1$ elements of the array beta. Beta (1) is set to zero.

## 6. Discussion of Numerical Properties

The behaviour of the reorthogonalization process is far from obvious. A detailed error analysis tends to obscure the essential simplicity of the underlying mechanism and we content ourselves with an exposition of the latter. For convenience it will be assumed that $\|c\|_{2}=1$.

We proceed by induction. Let us assume that on a computer with t-digit mantissa $X_{1}, x_{2}, \ldots, x_{r}$ have been determined and satisfy

$$
\begin{equation*}
\gamma_{i} x_{i+1}=C x_{i}-\alpha_{i} x_{i}-\beta_{i} x_{i-1}+O\left(2^{-t}\right) \quad i=1, \ldots, r-1 \tag{i}
\end{equation*}
$$

and

$$
\begin{equation*}
X_{i}^{T} X_{j}=\delta_{i j}+0\left(2^{-t}\right), \quad i, j<r \tag{ii}
\end{equation*}
$$

In other words we assume that the $\mathrm{x}_{\mathrm{i}}$ produced by the reorthogonalization technique are orthogonal to working accuracy.

In the next step $\mathrm{y}_{\mathrm{r}+1}$ is first determined and the computed vector satisfies the relation

$$
\begin{equation*}
y_{r+1}=C x_{r}-\alpha_{r} x_{r}-\beta_{r} x_{r-1}+O\left(2^{-t}\right) \tag{iii}
\end{equation*}
$$

If a great deal of cancellation takes place $y_{r+1}$ will not be accurately orthogonal to $x_{r}$ and $x_{r I}$ but the determination of $\alpha_{r}$ and $\beta_{r}$ ensures that $y_{r+1}^{T} x_{r}, y_{r+1}^{T} X_{r-1}=O\left(2^{-t}\right)$ (N.B. This will only imply accurate orthogonality if $\left\|y_{r+1}\right\|_{2}$ were of order unity; if $\left\|y_{r+1}\right\|_{2}$ is small, $y_{r+1}^{T} x_{r}$ can be of order $2^{-t}$ without $y_{r+1}$ being orthogonal to $X_{r}$ to working accuracy.) We now show that $y_{r+I}^{T} x_{i}=O\left(2^{-t}\right)$ for all earlier $x_{i}$. In fact we have

$$
\begin{align*}
y_{r+1}^{T} x_{i} & =x_{i}^{T}\left(C x_{r}-a_{r} x_{r}-\beta_{r} x_{r-1}\right)+0\left(2^{-t}\right)  \tag{iv}\\
& =x_{r}^{T} C x_{i}+0\left(2^{-t}\right) \\
& =x_{r}^{T}\left(x_{i+1}+\alpha_{i} x_{i}+\beta_{i} x_{i-1}\right)+0\left(2^{-t}\right) \\
& =O\left(2^{-t}\right) .
\end{align*}
$$

The essential point is that the inner-products of $y_{r+1}$ with respect to $x_{1}, \ldots, x_{r}$ are all negligible. If in particular $\left\|y_{r+1}\right\|_{2}$ is of the order of unity $\mathrm{y}_{\mathrm{r}+1}$ will already be accurately orthogonal to $x_{1} 0$ and reorthogonalization will be unnecessary. In any case we may write

$$
\begin{equation*}
y_{r+1}=\alpha_{z}+\eta_{I} x_{1}+\eta_{2} x_{2}+\cdots+\eta_{r} x_{r}, \quad\left|\eta_{i}\right|=0\left(2^{-t}\right) \tag{v}
\end{equation*}
$$

and

$$
\begin{equation*}
\alpha=\left\|y_{r+1}\right\|+0\left(2^{-t}\right)=\gamma_{r+1}+0\left(2^{-t}\right) \tag{vi}
\end{equation*}
$$

The vector $y_{r+1}$ is now multiplied by $P_{1}, P_{2}, \ldots, P_{r}$ successively and the resulting vector is used to determine $P_{r+1}$. From the derivation of the previous $P_{i}$ it is evident that

$$
\begin{gather*}
P_{r+1} P_{r} \ldots P_{1} y_{r+1}=\alpha e_{r+1} \eta_{1} e_{1}+\cdots \eta_{r} e_{r} 0\left(2^{-t}\right) \\
=\gamma_{r+1} e_{r+1}+0\left(2^{-t}\right) \tag{vii}
\end{gather*}
$$

where $e_{i}$ denotes the i-th column of $I$. Notice on the right-hand side of (vii), $\quad \gamma_{r+1}{ }^{e_{r+1}}$ may not necessarily be much larger than the term
denoted by $O\left(2^{-t}\right)$. If a great deal of cancellation took place when $\mathrm{y}_{\mathrm{r}+\mathrm{I}}$ was computed then $\mathrm{y}_{\mathrm{r}+\mathrm{I}}$ will be correspondingly small. However, independent of the size of ${ }^{y}{ }_{r+1}$ we have

$$
\begin{equation*}
y_{r+1}=\gamma_{r+1} P_{1} P_{2} \ldots P_{r+1} e_{r+1}+O\left(2^{-t}\right) \tag{viii}
\end{equation*}
$$

and substituting this in (iii) we have

Hence taking $x_{r+1}=P_{1} P_{2} \ldots P_{r+1} e_{r+1}$, Equation (ix) becomes

$$
\begin{equation*}
\gamma_{r+1} x_{r+1}=C x_{r}-\alpha_{r} x_{r}-\beta_{r} x_{r-1}+0\left(2^{-t}\right) \tag{x}
\end{equation*}
$$

and since earlier $x_{i}$ have been determined via the relation

$$
\begin{equation*}
\mathbf{x}_{i}=P_{1} P_{2} \ldots P_{i} e_{i} \tag{xi}
\end{equation*}
$$

it is clear that $x_{r+1}$ is orthogonal to all earlier $x_{i}$ to working accuracy. If there had been exact computation throughout, $\mathrm{x}_{\mathrm{r}+1}$ would have been $y_{r+1} /\left\|y_{r+1}\right\|_{2}=z_{r+1}$ (say). If cancellation has taken place and $\left\|y_{r+1}\right\|=2^{-k}$ (say), then (vii) shows that we can expect the computed $\left\|x_{r+1}-z_{r+1}\right\|$ to be of the order of $2^{k-t}$. Hence as $k$ becomes larger and approaches $t, x_{r+1}$ increasingly diverges from ${ }^{2} r+1$. However, since $\gamma_{r+1}=2^{-k}$ the replacement of $y_{r+1}$, i.e., $\gamma_{r+1} z_{r+1}$ by $\gamma_{r+1}{ }^{x}{ }_{r+1}$ on the left of Equation (iii) is merely a change of order $2^{-t}$. compare the computed $\beta_{i}$ and $\gamma_{i} \ldots$. We have

$$
\begin{aligned}
\beta_{i}=x_{i-1}^{T} C x_{i}+O\left(2^{-t}\right) & =x_{i}^{T}\left(C x_{i-1}\right)+0\left(2^{-t}\right) \\
& =x_{i}^{T}\left(\alpha_{i-1} x_{i-1}+\beta_{i-1} x_{i-2}+\gamma_{i} x_{i}+0\left(2^{-t}\right)\right)+0\left(2^{-t}\right) \\
& =\gamma_{i}+0\left(2^{-t}\right) .
\end{aligned}
$$

(Without the normalization of $C$ we have $\beta_{i}=\gamma_{i}+O\left(2^{-t}\|C\|_{2}\right)$.) Since $\beta_{i}$ and $y_{1}$ are floating point numbers, the number of figures agreeing in the mantissa depends on the degree of cancellation. But it is clear that if we replace $\beta_{i}$ by $\gamma_{i}$ we still have, as before,

$$
\begin{equation*}
\gamma_{i+1} x_{i+1}=C x_{i}-\alpha_{i} x_{i}-\gamma_{i} x_{i-1}+O\left(2^{-t}\right) \tag{xiii}
\end{equation*}
$$

and hence we can take the derived tridiagonal matrix to be symmetric. In the case when $y_{r+1}$ is zero (or is considered to be negligible) we can clearly take $x_{r+1}$ to be $P_{1} P_{2} . . . P_{r} e_{r+1}$ and we have

$$
0 \cdot x_{r+1}=\theta=y_{r+1}=C x_{r}-\alpha_{r} x_{r}-\beta_{r} x_{r-1}+O\left(2^{-t}\right)
$$

In this case $\beta_{r+1}$ will also turn out to be negligible to working accuracy.

An error analysis of the symmetric Lanczos process with Schmidt reorthogonalization has been given by Paige [2].
7. Test Results

|  | 10 | 2 | 3 | 1 | 1 |  | 12 | 1 | -1 | 2 | 1 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 2 | 12 | 12 |  | 1 |  | 1 | 14 | 1 | -1 | 1 |
| $A=$ | 3 | 1 | 11 | 1 | -1 | $B=$ | -1 | 1 | 16 | -1 | 1 |
|  | 1 | 2 | 1 | 9 | 1 | . | 2 | $-1$ | -1 | 12 | -1 |
|  | 1 | 1 | -1 | 1 | 15 |  | 1 | 1 | 1 | -1 | 11 |

were used. $A$ and $B$ are of full width, so $m a=m b=4$. On an IBM System 360 model 67 computer using floating point arithmetic with
a 14 hexadecimal digit fraction, the following results were obtained (Although not necessary, the elements of $u$ were initially zeroed.):

| $\alpha$ | 8 |
| :---: | :---: |
| $0.833333333333333_{10}+00$ | $0.000000000000000_{10}+00$ |
| $0.726877633595368_{10}+00$ | $-0.288543403757058_{10}+00$ |
| $0.116237235917115_{10}+01$ | $-0.217837154467399_{10}+00$ |
| $0 .{ }_{105692992323769_{10}+01}$ | $0.302923727655704_{10}+00$ |
| $0.862433487300640_{10}+00$ | $0.219669706658649_{10}+00$ |


| $0.000000000000000_{10}+\infty$ | $0.000000000000000_{10}+00$ | $0.000000000000000{ }_{10}+00$ |
| :---: | :---: | :---: |
| $0.109306814617572_{10}+00$ | $0.378822780886040_{10}+00$ | $0.000000000000000_{10}+00$ |
| $0.177 .357463219424{ }_{10} 001$ | $0.271519129954777{ }_{10}+00$ | $0.222348416368022_{10}+00$ |
| $671791653895_{10}+00$ | $-0.334341703878002_{10}-01$ | $-0.108503313670229_{10}+00$ |
| 0. $965095600469936_{10}-01$ | $-0.163232422022066{ }_{10}-01$ | $0.188837775101004_{10}+00$ |


$0.000000000000000_{10}+00$
$0.000000000000000{ }_{10}+00$
$0.000000000000000_{10}+00$
$0.000000000000000_{10}+\infty 0$
$-0.439339413317297_{10}+00$

The resulting tridiagonal system was solved using the procedure tiq12 After the vectors of the tridiagonal system were transformed according to equations (31) and (32), the final results were essentially the same as those reported in [1] for the above matrices.
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## Keywords

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